

## FINITE DIMENSIONALITY OF THE ATTRACTOR FOR A SEMILINEAR WAVE EQUATION WITH NONLINEAR BOUNDARY DISSIPATION

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**Abstract.** This paper deals with the long-time behavior of a semilinear wave equation with nonlinear boundary dissipation. By estimating energy function, we prove that the fractal dimension of the global attractor for the dynamical system is finite.

### 1. INTRODUCTION

Let  $\Omega \subset R^3$  be a bounded connected open set with a smooth boundary  $\Gamma$ . The exterior normal on  $\Gamma$  is denoted by  $\nu$ . We consider the following equation

$$u_{tt} - \Delta u + \varepsilon u_t + f(u) = 0 \quad \text{in } Q = [0, \infty) \times \Omega \quad (1.1)$$

subject to the boundary condition

$$\partial_\nu u + u = -g(u_t) \quad \text{in } \Sigma = [0, \infty) \times \Gamma \quad (1.2)$$

and the initial condition:

$$u(0) = u_0 \quad \text{and} \quad u_t(0) = u_1.$$

Here  $f$  and  $g$  are nonlinear functions subject to the following assumption,  $0 \leq \varepsilon \ll 1$  is a small parameter.

**Assumption 1.** (H1)  $f \in C^2(R)$  such that  $|f''(s)| \leq c(1 + |s|^{1-\delta})$  for all  $s \in R$  and for some  $c > 0$  and  $\delta > 0$ .

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(H2)  $\lim_{|s| \rightarrow \infty} s^{-1}f(s) > -\lambda$ , where  $\lambda$  is the best constant in the Poincare type inequality:  $\int_{\Omega} |\nabla u|^2 dx + \int_{\Gamma} |u|^2 dS \geq \lambda \int_{\Omega} |u|^2 dx$ .

(g)  $g \in C^1(R)$  is an increasing function,  $g(0) = 0$  and there exist two positive constants  $m_1$  and  $m_2$  such that  $m_1 \leq g'(s) \leq m_2$  for all  $|s| \geq 1$ .

This problem of long time behavior of solutions to (1.1)-(1.2) was the subject of a recent paper [2]. In [2] it was shown, in particular, that all finite energy solutions are attracted by a global compact attractor, whose structure is determined by the unstable manifolds emanating from stationary solutions.

Fereisel[4] dealt either with one-dimensional wave equations or with equations subjected to severe restrictions (linear bound) imposed on the semilinear terms [8]. Most recently, there has been a renewed interest in this problem and positive results for semilinear wave equations in dimension higher than one with interior dissipation have been established by Prazak [7].

Igor Chueshov [1] has shown under the Assumption 1.1 the fractal dimension of the global attractor of the (1.1) is finite when  $\varepsilon = 0$ . His methods can be extended to  $0 < \varepsilon \leq 1$ . In this paper we will use the methods to study the (1.1). Our main results read as follows:

**Theorem 1.** Under the Assumption 1.1 with  $g'(0) > 0$ , the fractal dimension of the global attractor of the dynamical system (1.1) is finite when  $0 < \varepsilon < \varepsilon_0$  where  $\varepsilon_0$  is a constant.

In order to describe the decay rates to equilibrium, we introduce a concave, strictly increasing, continuous function  $h : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  which captures the behavior of  $g(s)$  at the origin possessing the properties

$$h(0) = 0 \quad \text{and} \quad s^2 + g^2(s) \leq h(sg(s)) \quad \text{for } |s| \leq 1.$$

Given function  $h$  we define

$$H_0(s) = h\left(\frac{s}{c_3}\right), \quad G_0(s) = c_1(I + H_0)^{-1}(c_2s), \quad Q(s) = s - (I + G_0)^{-1}(s)$$

where  $c_1$  and  $c_2$  are positive constants depending on  $m_1$ ,  $m_2$  and  $c_3$  is proportional to the measure of  $\Gamma \times (0, T)$ .

## 2. PROOF OF THEOREM 1

We recall that the fractal dimension  $\dim_f M$  of a compact set  $M$  can be defined by the formula [3]

$$\dim_f M = \limsup_{\varepsilon \rightarrow 0} \frac{\ln N(M, \varepsilon)}{\ln(1/\varepsilon)},$$

where  $N(M, \varepsilon)$  is the minimal number of closed sets of diameter  $2\varepsilon$  which cover  $M$ .

Throughout the paper we shall use the energy functional

$$\begin{aligned} E(w(t)) &= \frac{1}{2} \int_{\Omega} |\nabla w(t)|^2 + \frac{1}{2} \int_{\Omega} |w_t(t)|^2 + \frac{1}{2} \int_{\Gamma} |w(t)|^2 \\ &\equiv \frac{1}{2} \| (w(t), w_t(t)) \|_{\mathcal{F}}^2, \end{aligned} \quad (2.1)$$

which corresponds to a linear version ( $f \equiv 0$  and  $g \equiv 0, \varepsilon \equiv 0$ ) of problem (1.1) and (1.2),  $\mathcal{F} \equiv H^1(\Omega) \times L^2(\Omega)$  which is often referred to as the finite energy space.

The key inequality in our considerations is the following inequality.

**Lemma 1.** Let  $u(t)$  and  $v(t)$  be two solutions to (1.1) possessing the properties

$$\| (u(t), u_t(t)) \|_{\mathcal{F}}^2 \leq R \quad \text{and} \quad \| (v(t), v_t(t)) \|_{\mathcal{F}}^2 \leq R \quad \text{for all } t \geq 0,$$

with some constant  $R > 0$ . Denote  $z(t) \equiv u(t) - v(t)$ . Then, under the Assumption 1 and  $g'(0) > 0$ , there exist positive constants  $C_1, C_2$  and  $\beta$  such that

$$E(z(t)) \leq C_1 e^{-\beta t} E(z(0)) + C_2 \int_0^t e^{-\beta(t-s)} \| z(s) \|_{L^2(\Omega)}^2 ds \quad (2.2)$$

for all  $t \geq 0$ .

**Remark.** Since the global attractor (we write  $\mathcal{A}$ ) is an invariant set we know that the solutions  $(u(t), u_t(t))$  and  $(v(t), v_t(t))$  corresponding to the initial data  $(u_0, u_1), (v_0, v_1) \in \mathcal{A}$  will stay in  $\mathcal{A}$ . Moreover, (2.2) implies that

$$E(z(t)) \leq C_1 e^{-\beta t} E(z(0)) + \frac{C_2}{\beta} \max_{s \in [0, t]} \| z(s) \|_{L^2(\Omega)}^2 ds \quad \text{for all } t \geq 0.$$

This inequality will be crucial in the proof of the theorem.

To prove Lemma 1 we need the following observability inequality which is valid without the restriction  $g'(0) > 0$ .

**Lemma 2.** Let  $T > T_0 \equiv 4(r + \frac{1}{\sqrt{\lambda}})$ , where  $r$  is the radius of a minimal ball in  $\mathbb{R}^n$  containing  $\Omega$  and  $\lambda$  is the constant from (H2). Assume that two solutions  $u(t)$  and  $v(t)$  to problem (1.1) possess the property

$$\| (u(t), u_t(t)) \|_{\mathcal{F}} \leq R \quad \text{and} \quad \| (v(t), v_t(t)) \|_{\mathcal{F}} \leq R \quad \text{for all } t \in [0, T]$$

with some constant  $R > 0$ . Then, under the Assumption 1 there exist positive constants  $C_1(T)$  and  $C_2(R, T)$  such that for  $z(t) \equiv u(t) - v(t)$  we have the

relations

$$\begin{aligned} E(z(T)) + \int_0^T E(z(t))dt &\leq C_2 \int_0^T \|z\|_{L^2(\Omega)}^2 dt \\ &+ C_1(I + H_0) \left( \int_0^T \langle g(u_t(t)) - g(v_t(t)), z_t(t) \rangle dt \right), \end{aligned} \quad (2.3)$$

where  $H_0(s) = Th(s/T)$  and  $h(s)$  is defined in the beginning. If  $g'(0) > 0$ , then (2.3) hold with  $H_0(s) \equiv s$  and

$$E(z(T)) \leq C_1(E(z(0)) - E(z(T))) + C_2 \int_0^T \|z\|_{L^2(\Omega)}^2 dt. \quad (2.4)$$

*Proof.* We will denote different constant by  $C_i$  pointing out their dependence on the parameters when it becomes important.

Our starting point is the following identity:

$$\begin{aligned} \frac{1}{2} \int_{Q_T} (|z_t|^2 + |\nabla z|^2) &= - \int_{Q_T} (f(u) - f(v))(\tilde{h} \cdot \nabla z + z) \\ &- \int_{\Omega} z_t(\tilde{h} \cdot \nabla z + z)|_0^T - \int_{Q_T} \varepsilon z_t(\tilde{h} \cdot \nabla z + z) + \int_{\Sigma_T} \left[ \partial_\nu z(\tilde{h} \cdot \nabla z + z) \right. \\ &\left. + \frac{1}{2}((\tilde{h} \cdot \nu)(|z_t|^2 - |\nabla z|^2)) \right], \end{aligned} \quad (2.5)$$

where  $\tilde{h} = x - x_0$  for some  $x_0 \in R^3$  and  $Q_T = (0, T) \times \Omega$  and  $\Sigma_T = (0, T) \times \Gamma$ . Here  $T$  is a positive constant which will be determined later in the proof. We choose  $x_0 \in R^3$  such that  $\sup_{x \in \Omega} |\tilde{h}(x)| = r$ , where  $r$  is the radius of a minimal ball in  $R^3$  containing  $\Omega$ .

In the following we will transform (2.5) to an energy inequality.

At first we estimate the second term in (2.5). By

$$f(u) - f(v) = (u - v) \int_0^1 f'(su + (1-s)v)ds = z \int_0^1 f'(su + (1-s)v)ds,$$

the Hölder's inequality, and the embedding  $H^s(\Omega) \subset L^p(\Omega)$  for  $s = \frac{3}{2} - \frac{3}{p}$  and  $p \geq 2$ , we can get that

$$\|f(u(t)) - f(v(t))\|_{L^2(\Omega)} \leq C(R, \Omega, c, \delta) \cdot \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}, \quad t \in [0, T], \quad (2.6)$$

for some  $\delta > 0$  and

$$\begin{aligned} &- \int_{\Omega} (f(u(t)) - f(v(t)))\tilde{h} \cdot \nabla z(t) \\ &\leq C(\varepsilon)r^2 \|f(u(t)) - f(v(t))\|_{L^2(\Omega)}^2 + \varepsilon \|\nabla z(t)\|_{L^2(\Omega)}, \end{aligned}$$

for any  $\varepsilon > 0$ . By Hölder's inequality, and the embedding  $H^s(\Omega) \subset L^p(\Omega)$  for  $s = \frac{3}{2} - \frac{3}{p}$  and  $p \geq 2$ ,

$$- \int_{\Omega} (f(u(t)) - f(v(t))) \tilde{h} \cdot z \leq C(R, \Omega, c, \delta) \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}, \quad t \in [0, T],$$

So we can conclude

$$- \int_{Q_T} (f(u(t)) - f(v(t))) \tilde{h} \cdot z \leq C \int_0^T \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)} dt.$$

It is easy to see that  $|\int_{\Omega} z_t (\tilde{h} \cdot \nabla z + z)| \leq (r + \frac{1}{\sqrt{\lambda}}) E(z(t))$ .

Hence, according to (2.5) and the definition of the energy functional (2.2) yield

$$\begin{aligned} \int_0^T E(z(t)) dt &\leq C_1(r) \left( \|z_t\|_{L^2(\Sigma_T)}^2 + \|\nabla z\|_{L^2(\Sigma_T)}^2 + \|z\|_{L^2(\Sigma_T)}^2 \right) \\ &+ \left( r + \frac{1}{\sqrt{\lambda}} \right) [E(z(T)) + E(z(0))] + \left( r + \frac{1}{\sqrt{\lambda}} \right) \varepsilon \int_0^T E(z(t)) dt \\ &+ C_2(r, R, \varepsilon) \int_0^T \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 dt + \int_0^T \|\nabla z\|_{L^2(\Omega)}^2 dt. \end{aligned} \quad (2.7)$$

Multiplying the differential equation for  $z$  by  $z_t$  and integrating by parts result in the following energy identity

$$\begin{aligned} E(z(t)) + \int_s^t \langle g(u_t(\tau)) - g(v_t(\tau)), z_t(\tau) \rangle d\tau + \varepsilon \int_s^t (z_t(\tau), z_t(\tau)) d\tau \\ + \int_s^t \langle f(u(\tau)) - f(v(\tau)), z_t(\tau) \rangle d\tau = E(z(s)) \quad \text{for } 0 \leq s \leq t. \end{aligned} \quad (2.8)$$

Since the second term and the third on the left side are non-negative, we can also write

$$E(z(t)) \leq E(z(s)) - \int_s^t \langle f(u(\tau)) - f(v(\tau)), z_t(\tau) \rangle d\tau \quad \text{for } 0 \leq s \leq t. \quad (2.9)$$

Therefore using (2.6), the embedding  $H^1(\Omega) \subset L^6(\Omega) \subset L^{\frac{6}{1+\delta}}(\Omega)$  and Gronwall's lemma we obtain from (2.9) that

$$E(z(t)) \leq E(z(s)) \cdot e^{a_R(t-s)} \quad \text{for all } 0 \leq s \leq t, \quad (2.10)$$

where the constant  $a_R > 0$  depends on  $\Omega$ ,  $c$  and  $\delta$ . By (g), without loss of generality, we can suppose that  $0 \leq g'(s) \leq m_2$  for all  $s \in R$ . Therefore we

derive from (2.9),(2.8) and (g) we have the following result

$$\begin{aligned} E(z(s)) &\leq E(z(t)) + m_2 \int_s^t \int_{\Gamma} |z_t(\tau)|^2 d\tau \\ &\quad + \int_s^t (2\varepsilon \|z_t(\tau)\|_{L^2(\Omega)}^2 + C(\varepsilon, R) \|z(\tau)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2) d\tau, \end{aligned} \quad (2.11)$$

for any  $\varepsilon > 0$  and for all  $0 \leq s \leq t$ . In a same way from (2.8) we have

$$\begin{aligned} \int_s^t \langle g(u_t(\tau)) - g(v_t(\tau)), z_t(\tau) \rangle d\tau &\leq E(z(s)) - E(z(t)) \\ &\quad + \int_s^t (2\varepsilon \|z_t(\tau)\|_{L^2(\Omega)}^2 + C(\varepsilon, R) \|z(\tau)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2) d\tau, \end{aligned} \quad (2.12)$$

for any  $\varepsilon > 0$  and for all  $0 \leq s \leq t$ .

The estimate of the tangential derivative  $\nabla_{\tau} z$  on  $\Sigma_T$  is given in Lemma 7.2 in Lasiecka and Irigiani (1992) below [6]. For  $0 < \alpha < T/2$  and  $\eta \in (0, 1/2)$  there exists a constant  $C = C(\alpha, \eta, T, \Omega)$  such that

$$\begin{aligned} \int_{\alpha}^{T-\alpha} \int_{\Gamma} |\nabla_{\tau} z(t)|^2 dt &\leq C \left( \|\partial_{\nu} z\|_{L^2(\Sigma_T)}^2 + \|z_t\|_{L^2(\Sigma_T)}^2 + \|z\|_{H^{\frac{1}{2+\eta}}(Q_T)}^2 \right. \\ &\quad \left. + \|f(u) - f(v)\|_{H^{-\frac{1}{2+\eta}}(Q_T)}^2 \right). \end{aligned}$$

By (2.6) the last term on the right-hand side can be estimated in the following

$$\begin{aligned} &\|f(u) - f(v)\|_{H^{-\frac{1}{2+\eta}}(Q_T)}^2 \\ &\leq C \|f(u) - f(v)\|_{L^2(Q_T)}^2 \leq C(R) \int_0^T \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 dt \end{aligned} \quad (2.13)$$

We derive from (2.11)

$$\begin{aligned} E(z(\alpha)) + E(z(T - \alpha)) &\leq 2E(z(T)) + 2m_2 \|z_t\|_{L^2(\Sigma_T)}^2 + 4\varepsilon \|z_t\|_{L^2(Q_T)}^2 \\ &\quad + C(\varepsilon, R) \int_0^T \|z(\tau)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 d\tau \end{aligned} \quad (2.14)$$

and

$$\begin{aligned} E(z(\alpha)) + E(z(T - \alpha)) &\leq 2\alpha E(z(T)) + 2\alpha m_2 \|z_t\|_{L^2(\Sigma_T)}^2 \\ &\quad + 4\alpha\varepsilon \|z_t\|_{L^2(Q_T)}^2 + \alpha C(\varepsilon, R) \int_0^T \|z(\tau)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 d\tau \end{aligned} \quad (2.15)$$

for any  $\alpha \in (0, T/2)$ . Therefore using estimate (2.7) for the interval  $(\alpha, T - \alpha)$  and also relation (2.13) we obtain (after redefining  $\varepsilon$ ) and setting  $c_0 =$

$$2\left(\alpha + r + \frac{1}{\lambda}\right),$$

$$\begin{aligned} & \int_0^T E(z(t))dt \leq C_1\{\|z_t\|_{L^2(\Sigma_T)}^2 + \|\partial_\nu z\|_{L^2(\Sigma_T)}^2 + \|z\|_{L^2(\Sigma_T)}^2\} \\ & + C_2(\varepsilon, R) \left\{ \int_0^T \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 dt + \|z(t)\|_{H^{\frac{2}{2+\eta}}(Q_T)}^2 \right\} \\ & + \varepsilon \int_0^T (\|z_t(t)\|_{L^2(\Omega)}^2 + \|\nabla z(t)\|_{L^2(\Omega)}^2)dt + c_0 E(z(T)). \end{aligned} \quad (2.16)$$

Integrating now inequality (2.9) with  $t = T$  with respect to  $s$  over the interval  $[0, T]$  and using (2.6) yields

$$TE(z(T)) \leq \int_0^T E(z(s))ds + \varepsilon \|z_t\|_{L^2(Q_T)}^2 + C(R, T, \varepsilon) \int_0^T \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 dt.$$

Hence, choosing  $\varepsilon$  small enough, from (2.16) we obtain that

$$\begin{aligned} & TE(z(T)) + \frac{1}{2} \int_0^T E(z(t))dt \\ & \leq C_1 \left( \|z_t\|_{L^2(\Sigma_T)}^2 + \|\partial_\nu z\|_{L^2(\Sigma_T)}^2 + \|z\|_{L^2(\Sigma_T)}^2 \right) \\ & + 2c_0 E(z(T)) + C_2 \left\{ \int_0^T \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 dt + \|z\|_{H^{\frac{1}{2+\eta}}(Q_T)}^2 \right\}. \end{aligned}$$

Therefore for any  $T > T_0 \equiv 4\left(r + \frac{1}{\sqrt{\lambda}}\right)$  we can choose appropriate  $\alpha$  and obtain an estimate of the form

$$\begin{aligned} & TE(z(T)) + \int_0^T E(z(t))dt \\ & \leq C_1(T) \left( \|z_t\|_{L^2(\Sigma_T)}^2 + \|\partial_\nu z\|_{L^2(\Sigma_T)}^2 + \|z\|_{L^2(\Sigma_T)}^2 \right) \\ & + C_2(R, T) \left\{ \int_0^T \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 dt + \|z\|_{H^{\frac{1}{2+\eta}}(Q_T)}^2 \right\}. \end{aligned} \quad (2.17)$$

For the boundary terms we can get the result from the proof of [1]. Since

$$\partial_\nu z + z = -g(u_t) + g(v_t) = -z_t \int_0^1 g'(su_t + (1-s)v_t)ds,$$

using the inequality  $0 \leq g' \leq m_2$  we have

$$\|z_t\|_{L^2(\Sigma_T)}^2 + \|\partial_\nu z\|_{L^2(\Sigma_T)}^2 \leq (1 + 2m_2^2) \|z_t\|_{L^2(\Sigma_T)}^2 + 2 \|z\|_{L^2(\Sigma_T)}^2.$$

By Assumption (g) we can get

$$\begin{aligned} & \|z_t\|_{L^2(\Sigma_T)}^2 + \|\partial_\nu z\|_{L^2(\Sigma_T)}^2 \\ & \leq C_0 \int_0^T \langle g(z_t), z_t \rangle dt + C_1 \int_{\Sigma_T^*} |z_t|^2 ds dt + 2 \|z\|_{L^2(\Sigma_T)}^2, \end{aligned}$$

where  $\Sigma_T^* = \{(x, t) \in \Sigma_T : |z_t(x, t)| \leq 1\}$ . Since  $s^2 \leq h(g(s))$  for  $|s| \leq 1$ , using Jensen's inequality we obtain that

$$\begin{aligned} \|z_t\|_{L^2(\Sigma_T)}^2 + \|\partial_\nu z\|_{L^2(\Sigma_T)}^2 & \leq C \int_0^T (I + h)(\langle g(z_t), z_t \rangle) dt + 2 \|z\|_{L^2(\Sigma_T)}^2 \\ & \leq C \cdot (I + H_0) \left( \int_0^T \langle g(z_t), z_t \rangle dt \right) + 2 \|z\|_{L^2(\Sigma_T)}^2. \end{aligned} \quad (2.18)$$

Therefore (2.18) and (2.17) yield

$$\begin{aligned} & E(z(T)) + \int_0^T E(z(t)) dt \\ & \leq C_1(I + H_0) \left( \int_0^T \langle g(u_t(t)) - g(v_t(t)), z_t \rangle dt \right) + C_2 A(z), \end{aligned} \quad (2.19)$$

where  $C_1 = C_1(T)$ ,  $C_2 = C_2(R, T)$  and  $A(z)$  is an abbreviation for a collection of lower-order terms, that is,

$$A(z) = \int_0^T \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 dt + \|z\|_{H^{\frac{1}{2+\eta}}(Q_T)}^2 + \|z\|_{L^2(\Sigma_T)}^2.$$

From (2.12) and (2.19) we also have

$$E(z(T)) + \int_0^T E(z(t)) dt \leq C_1(I + H_0)(E(z(0)) - E(z(T))) + C_2 A(z), \quad (2.20)$$

In the following we will estimate  $A(z)$ . By the trace theorem we obtain that  $\|z(t)\|_{L^2(\Sigma_T)}^2 \leq C \|z(t)\|_{H^{\frac{1}{2+\eta}}(\Omega)}^2$ . The interpolation together with the

embedding  $H^{1-\frac{\delta}{2}}(\Omega) \subset L^{\frac{6}{1+\delta}}(\Omega)$  give us that

$$\begin{aligned} \|z\|_{H^{\frac{1}{2}+\eta}(Q_T)}^2 & \leq 2\varepsilon \|z\|_{H^1(Q_T)}^2 + C(\varepsilon) \|z\|_{L^2(Q_T)}^2 \\ & \leq c\varepsilon \int_0^T E(z(t)) dt + C(\varepsilon) \|z\|_{L^2(Q_T)}^2, \end{aligned}$$

$$\begin{aligned} \int_0^T \|z(t)\|_{L^{\frac{6}{1+\delta}}(\Omega)}^2 dt & \leq C \int_0^T \|z(t)\|_{H^{1-\frac{\delta}{2}}(\Omega)}^2 dt \\ & \leq \varepsilon \int_0^T E(z(t)) dt + C(\varepsilon) \|z\|_{L^2(Q_T)}^2. \end{aligned}$$



Thus we can draw the following estimate

$$A(z) \leq \varepsilon \int_0^T E(z(t))dt + C(\varepsilon) \|z\|_{L^2(Q_T)}^2, \tag{2.21}$$

with arbitrary  $\varepsilon > 0$ . Hence, (2.19) and (2.20) yield the desired relation.

In the case  $g'(0) > 0$  we have  $h(s) = c_0s$  and therefore (2.3) holds with  $H(s) = s$ . Similarly, (2.20) and (2.21) imply (2.4).  $\square$

**Proof of Lemma 2.** Under the hypotheses and  $g'(0) > 0$  it implies from (2.4) that

$$E(z(nT)) \leq \frac{C_1}{1 + C_1} E(z((n - 1)T)) + \frac{C_2}{1 + C_1} \int_{(n-1)T}^{nT} \|z\|_{L^2(\Omega)}^2 dt, \tag{2.22}$$

( $n = 1, 2, \dots$ ) for fixed  $T > T_0$ . We can prove by induction that

$$E(z(nT)) \leq \gamma^n E(z(0)) + \frac{C_2}{1 + C_1} \sum_{k=1}^n \gamma^{n-k} \int_{(k-1)T}^{kT} \|z\|_{L^2(\Omega)}^2 dt \tag{2.23}$$

for all positive integers  $n$ , where  $\gamma = C_1/(1 + C_1)$ . It follows from (2.10) that

$$E(z(t)) \leq E(z(nT)) \cdot e^{a_R T} \text{ for all } nT \leq (n - 1)T, \quad n = 0, 1, 2, \dots \tag{2.24}$$

Set  $\beta = \frac{1}{T} \ln \frac{1}{\gamma}$ . It is clear that  $\gamma^{n-k} \leq \frac{1}{\gamma^2} \exp -\beta(t - \tau)$  for  $t \leq (n + 1)T$  and  $\tau \geq (k - 1)T$ . Therefore (2.2) follows from (2.23) and (2.24).

In order to prove the theorem, we need the following theorem 2 [1]:

**Theorem 2.** Let  $X$  be a separable Hilbert space and  $A$  be a bounded closed set in  $X$ . Assume that there exists a mapping  $V : A \mapsto X$  such that

- (1)  $A \subset VA$ .
- (2)  $V$  is Lipschitz on  $A$ , i.e, there exists  $L > 0$  such that

$$\|Va_1 - Va_2\| \leq L \|a_1 - a_2\| \text{ for all } a_1, a_2 \in A.$$

- (3) There exist compact seminorms  $n_1(x)$  and  $n_2(x)$  on  $X$  such that

$$\|Va_1 - Va_2\| \leq \eta \|a_1 - a_2\| + K \cdot [n_1(a_1 - a_2) + n_2(Va_1 - Va_2)],$$

for all  $a_1, a_2 \in A$ , where  $0 < \eta < 1$  and  $K > 0$  are constants.

Then  $A$  is a compact set in  $X$  of a finite fractal dimension with the estimate for the following dimension

$$\dim_f A \leq d \cdot \ln \left( 1 + \frac{4(1 + L)}{\alpha} \right) \cdot \left[ \ln \frac{1}{\eta + \delta + \alpha \tilde{K}} \right]^{-1}.$$

In a special case when the seminorms  $n_i$  have the form  $n_i(v) = \|P_i v\|$ , where  $P_1$  and  $P_2$  are finite-dimensional orthoprojections, we have that

$$\dim_f A \leq (\dim P_1 + \dim P_2 \cdot \ln \left( 1 + \frac{8(1+L)\sqrt{2}K}{1-\eta} \right)) \cdot \left[ \ln \frac{2}{1+\eta} \right]^{-1}.$$

**Proof of Theorem 1.** Set the space  $X = \mathcal{F} \times H_1(Q_T)$  equipped with the norm

$$\|U\|_X = \|\nabla u_0\|_{L^2(\Omega)}^2 + \|u_0\|_{L^2(\Gamma)}^2 + \|u_1\|_{L^2(\Omega)}^2 + \int_0^T E(v(t))dt,$$

where  $U = (u_0, u_1, v)$ ,  $T > 0$  is a constant to be determined later. On the space  $X$  we define a seminorm

$$n_T(U) = \max_{0 \leq t \leq T} \|v(t)\|_{L^2(\Omega)},$$

then  $n_T(U)$  is a compact seminorm on  $X$  by the compactness of the imbedding [9]. Consider in the space  $X$  the set

$$\mathcal{A}_T = \{U \equiv (u_0, u_1, u(t) \text{ for } t \in [0, T]) : (u_0, u_1) \in \mathcal{A}\},$$

where  $u(t)$  is the solution to (1.1) with initial data  $u(0) = u_0$ ,  $u_t(0) = u_1$  and  $\mathcal{A}$  is an attractor. We define the operator  $V_T : \mathcal{A} \mapsto X$  by the formula

$$V_T : (u_0, u_1, u(t)) \mapsto (u(T), u_t(T), u(T+t)) = (s(T)(u_0, u_1), u(T+t)).$$

We shall show that all conditions of Theorem 2 are satisfied. For (1), this follows from the invariance property of the attractor  $\mathcal{A}$  which is equivalent to  $V_T \mathcal{A}_T = \mathcal{A}_T$ . As for (2),  $V_T$  is Lipschitz continuous on  $\mathcal{A}_T$ . In order to show this statement we will work with two solutions  $u(t)$  and  $v(t)$  to the problem (1.1). We set  $U_1 = (u_0, u_1, u(t))$ ,  $U_2 = (v_0, v_1, v(t))$  and  $z(t) = u(t) - v(t)$  and observe that

$$\begin{aligned} \frac{1}{2} \|U_1 - U_2\|_X^2 &= E(z(0)) + \int_0^T E(z(t))dt, \text{ and} \\ \frac{1}{2} \|V_T U_1 - V_T U_2\|_X^2 &= E(z(T)) + \int_T^{2T} E(z(t))dt, \end{aligned} \tag{2.25}$$

From (2.10), we can get

$$\int_T^{2T} E(z(s))ds \leq e^{a_R T} \int_0^T E(z(s))ds.$$

We have the Lipschitz property of  $V_T$  with  $L = e^{a_R T/2}$  when we combine the above inequality and formula (2.10).

Integrating the inequality in the Remark over the interval  $[T, 2T]$  obtain

$$\int_T^{2T} E(z(t))dt \leq C_1 e^{-\beta T} E(z(0)) + C_2 T \max_{0 \leq \tau \leq 2T} \|z(\tau)\|_{L^2(\Omega)}^2, \quad (2.26)$$

where  $C_1$  and  $C_2$  do not depend on  $T$ . Combining the last inequality and the inequality in Remark, we can get

$$E(z(T)) + \int_T^{2T} E(z(t))dt \leq C_1 e^{-\beta T} E(z(0)) + C_2 \max_{0 \leq \tau \leq 2T} \|z(\tau)\|_{L^2(\Omega)}^2. \quad (2.27)$$

Since

$$\max_{0 \leq \tau \leq 2T} \|z(\tau)\|_{L^2(\Omega)}^2 \leq \max_{0 \leq \tau \leq T} \|z(\tau)\|_{L^2(\Omega)}^2 + \max_{0 \leq \tau \leq T} \|z(\tau + T)\|_{L^2(\Omega)}^2,$$

accounting for the definitions of  $V_t$  and the norms in  $X$ , relation (2.27) can be written in the following form

$$\|V_T U_1 - V_T U_2\|_X^2 \leq \eta_T \|U_1 - U_2\|_X^2 + K \cdot [\eta_T (U_1 - U_2) + \eta_T (V_T U_1 - V_T U_2)]$$

for all  $U_1, U_2 \in \mathcal{A}_T$ , where  $\eta_T = C_1 e^{-\beta T}$ . We can select  $T$  large enough such that  $\eta_T < 1$ .

Hence, all the assumption of Theorem 2 are satisfied. It implies that  $\mathcal{A}_T$  is a compact set in  $X$  of finite fractal dimension. Let  $\mathcal{P} : X \rightarrow \mathcal{F}$  be the operator defined by the formula

$$\mathcal{P} : (u_0, U_1, v(t)) \rightarrow (u_0, u_1).$$

Since  $\mathcal{A} = \mathcal{P}\mathcal{A}_T$  and  $\mathcal{P}$  is obviously Lipschitz continuous, we have that

$$\dim_{frac}^{\mathcal{F}} \mathcal{A} = \dim_{frac}^X \mathcal{A}_T < \infty$$

where  $\dim_{frac}^Y$  stands for fractal dimension of a set in the space  $Y$ .

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